



Single transverse spin asymmetry (SSA) of W/Z bosons

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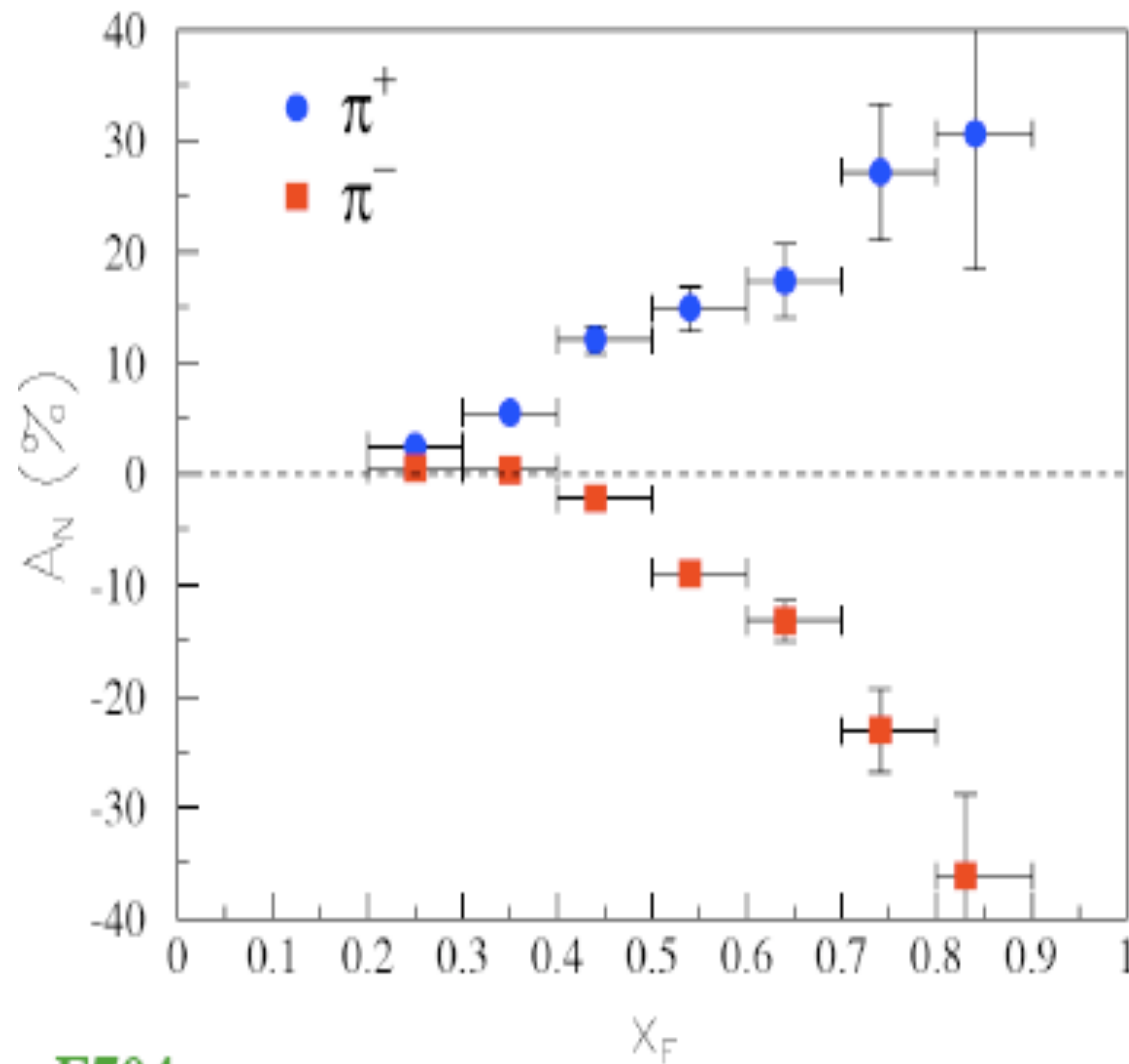
Berkeley Summer Program on Nuclear Spin Physics
Berkeley, CA, June 11, 2009

based on work with J. W. Qiu

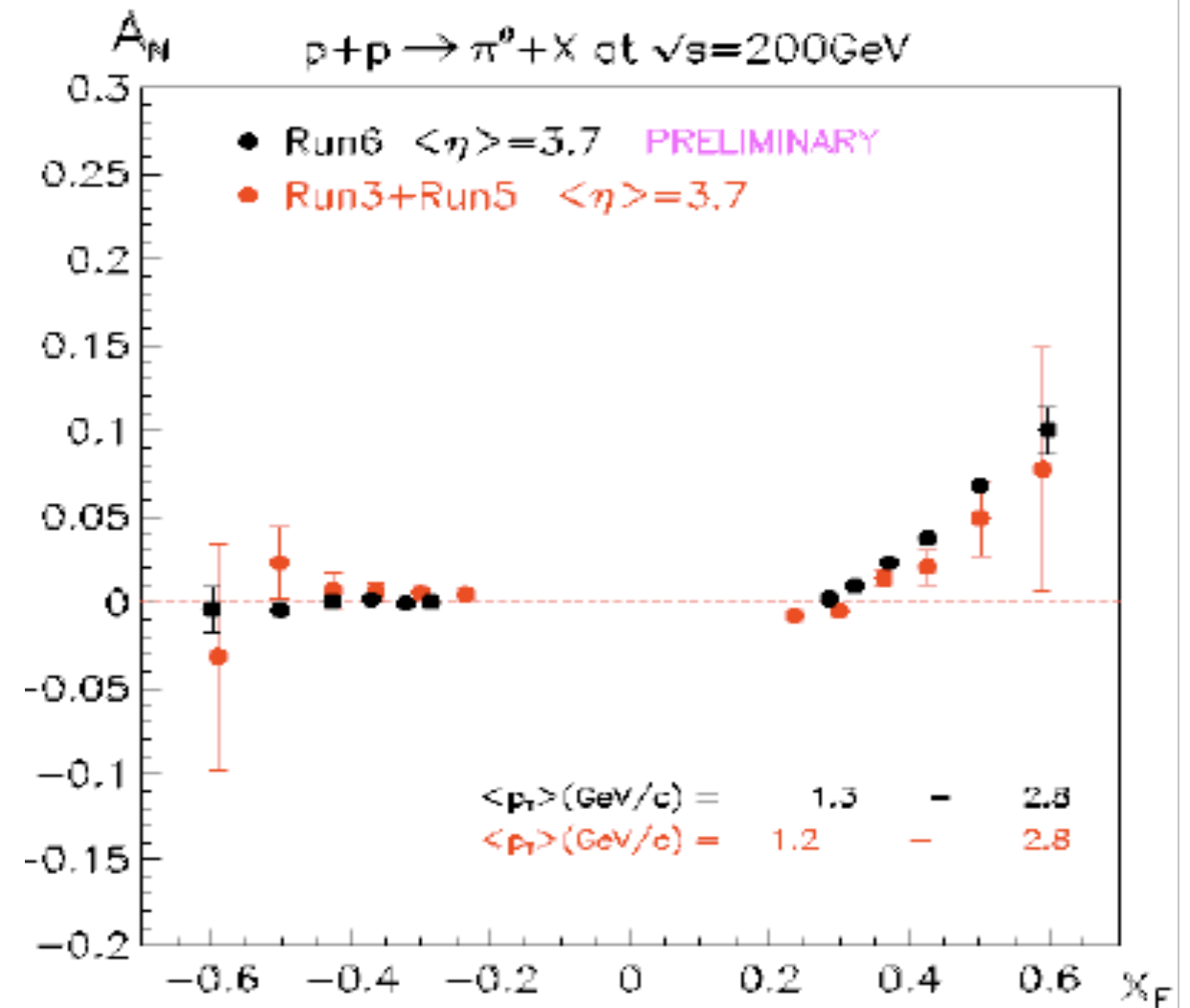
Experiment: Single Spin Asymmetries

- Fermilab E704, STAR, PHENIX, BRAHMS, COMPASS, HERMES, JLAB:

$$p^\uparrow p \rightarrow \pi X$$



E704

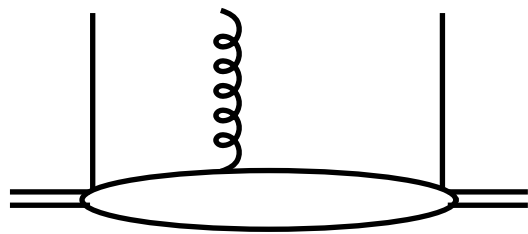


STAR (BRAHMS, too)

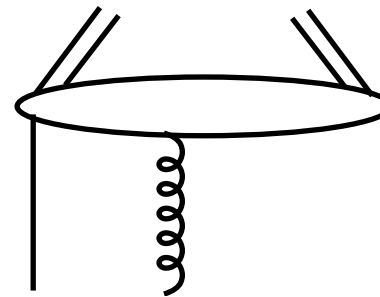
SSAs are observed in various experiments at different \sqrt{s}

Two mechanisms to generate SSA in QCD

- SSA is related to parton's transverse motion
- Collinear factorization approach:
 - Twist-3 three-parton correlation functions: Qiu-Sterman matrix element, ...
 - Twist-3 three-parton fragmentation functions:



Efremov-Teryaev 82, 84, Qiu-Sterman 91, 98, ...



Koike, 02, Zhou, Yuan, 09

- TMD approach: **T**ransverse **M**omentum **D**ependent distributions probe the parton's intrinsic transverse momentum
 - Sivers function: in Parton Distribution Function (PDF)
Sivers 90
 - Collins function: in Fragmentation Function (FF)
Collins 93

Relation between twist-3 and TMD approaches

- They apply in different kinematic domain:

- TMD approach: need TMD factorization, applies for the process with two observed momentum scales: DY at small q_T

$$Q_1 \gg Q_2 \left\{ \begin{array}{l} Q_1 \text{ necessary for pQCD factorization to have a chance} \\ Q_2 \text{ sensitive to parton's transverse momentum} \end{array} \right.$$

- Collinear factorization approach: more relevant for single scale hard process: inclusive pion production at pp collision

- They generate same results in the overlap region when they both apply:

- Twist-3 three-parton correlation in distribution \longleftrightarrow Sivers function

Ji, Qiu, Vogelsang, Yuan, 06, ...

- Twist-3 three-parton correlation in fragmentation \longleftrightarrow Collins function

Zhou, Yuan, 09



Major difference in these two approaches

- Collinear factorization approach:
 - All the twist-3 correlation functions (both in distribution and fragmentation side) are universal
- However, the TMD function in TMD approach MIGHT not be universal
 - Sivers function is NOT universal

Collins 02, Boer, Mulders, Pijlman, 03, Collins, Metz, 04, Kang, Qiu, 09, ...
 - Collins function is universal

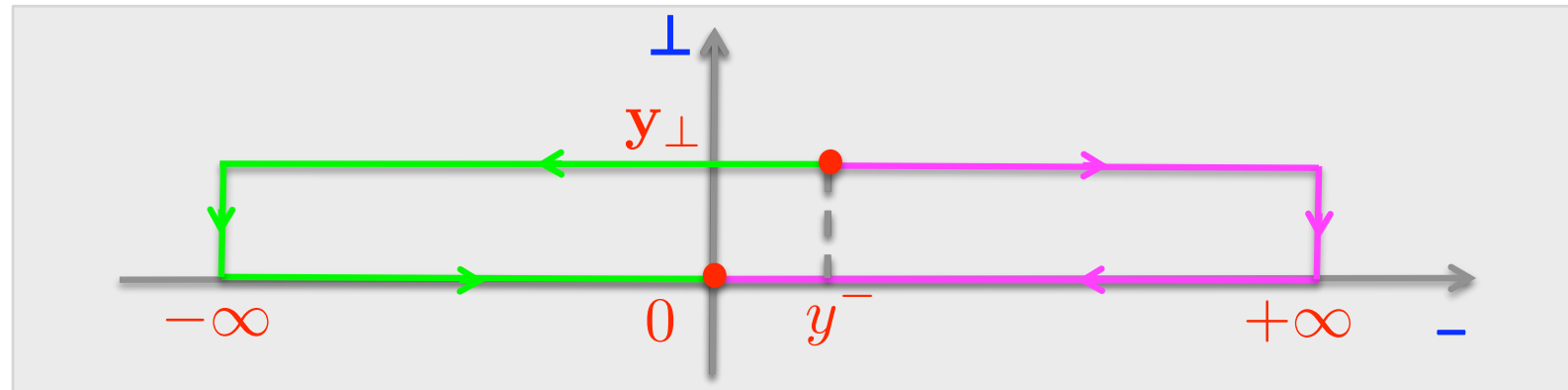
Metz 02, Collins, Metz, 04, Yuan, 08, Gamberg, Mukerjee, Mulders, 08, Meissner, Metz, 08, Zhou, Yuan, 09, ...

Non-universality of the Siverson function

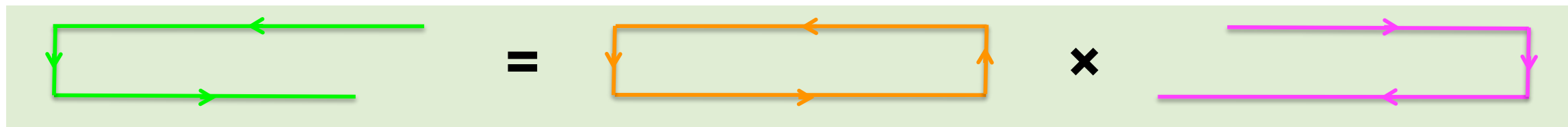
- Different gauge link for gauge-invariant TMD distribution in SIDIS and DY

$$f_{q/h^\uparrow}(x, \mathbf{k}_\perp, \vec{S}) = \int \frac{dy^- d^2 y_\perp}{(2\pi)^3} e^{ixp^+ y^- - i \mathbf{k}_\perp \cdot \mathbf{y}_\perp} \langle p, \vec{S} | \bar{\psi}(0^-, \mathbf{0}_\perp) \boxed{\text{Gauge link}} \frac{\gamma^+}{2} \psi(y^-, \mathbf{y}_\perp) | p, \vec{S} \rangle$$

- **SIDIS:** $\Phi_n^\dagger(\{+\infty, 0\}, \mathbf{0}_\perp) \Phi_{n_\perp}^\dagger(+\infty, \{\mathbf{y}_\perp, \mathbf{0}_\perp\}) \Phi_n(\{+\infty, y^-\}, \mathbf{y}_\perp)$
- **DY:** $\Phi_n^\dagger(\{-\infty, 0\}, \mathbf{0}_\perp) \Phi_{n_\perp}^\dagger(-\infty, \{\mathbf{y}_\perp, \mathbf{0}_\perp\}) \Phi_n(\{-\infty, y^-\}, \mathbf{y}_\perp)$



Wilson Loop $\sim \exp \left[-ig \int_{\Sigma} d\sigma^{\mu\nu} F_{\mu\nu} \right]$ Area is NOT zero



- For a fixed spin state:

$$f_{q/h^\uparrow}^{\text{SIDIS}}(x, \mathbf{k}_\perp, \vec{S}) \neq f_{q/h^\uparrow}^{\text{DY}}(x, \mathbf{k}_\perp, \vec{S})$$

Time-reversal modified universality of the Siverson function

- Relation between Siverson functions in SIDIS and DY

- From P and T invariance:

$$f_{q/h^\uparrow}^{\text{SIDIS}}(x, \mathbf{k}_\perp, \vec{S}) = f_{q/h^\uparrow}^{\text{DY}}(x, \mathbf{k}_\perp, -\vec{S})$$

- **Spin-averaged parton distribution function is universal**

- From definition:

$$f_{q/h^\uparrow}(x, \mathbf{k}_\perp, \vec{S}) \equiv f_{q/h}(x, k_\perp) + \frac{1}{2} \Delta^N f_{q/h^\uparrow}(x, k_\perp) \vec{S} \cdot \hat{p} \times \hat{\mathbf{k}}_\perp$$

- One can derive:

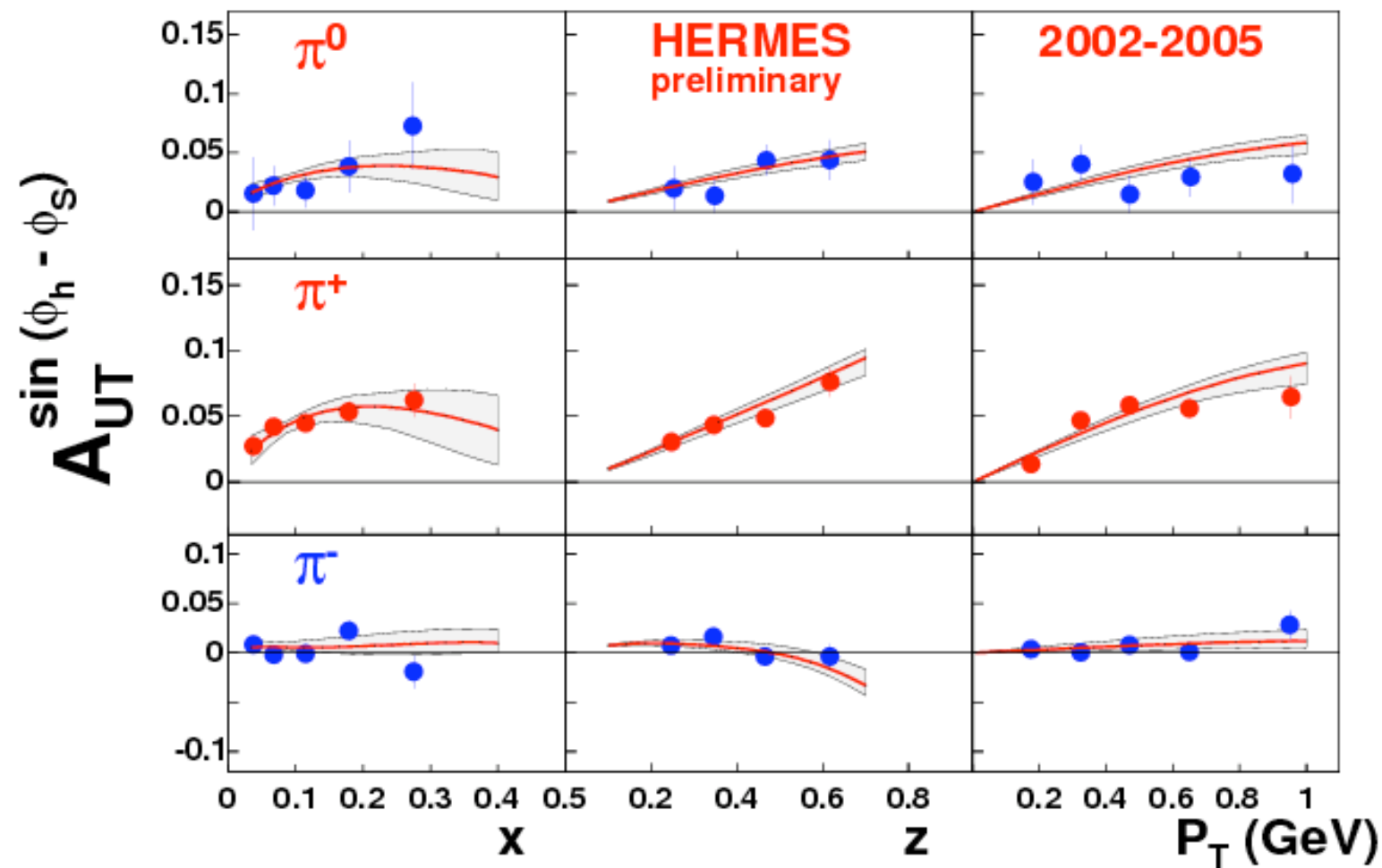
$$\Delta^N f_{q/h^\uparrow}^{\text{SIDIS}}(x, k_\perp) = -\Delta^N f_{q/h^\uparrow}^{\text{DY}}(x, k_\perp)$$

Most critical test for TMD approach to SSA

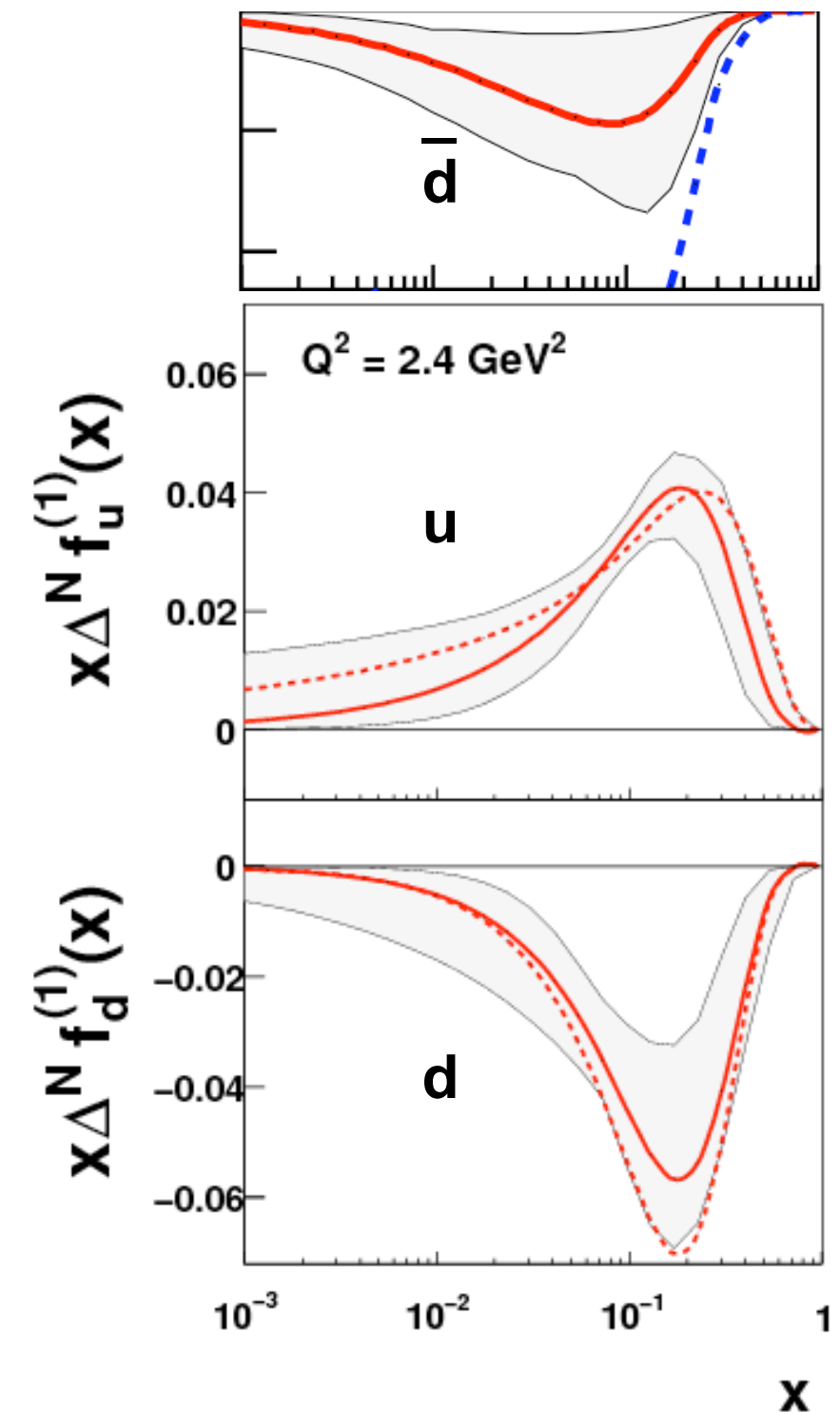
Sivers function from SIDIS

- Extract Sivers function from SIDIS

Anselmino, et.al., 2009



SIDIS		DY
Sivers _{u-quark} > 0	QCD →	Sivers _{u-quark} < 0
Sivers _{d-quark} < 0		Sivers _{d-quark} > 0

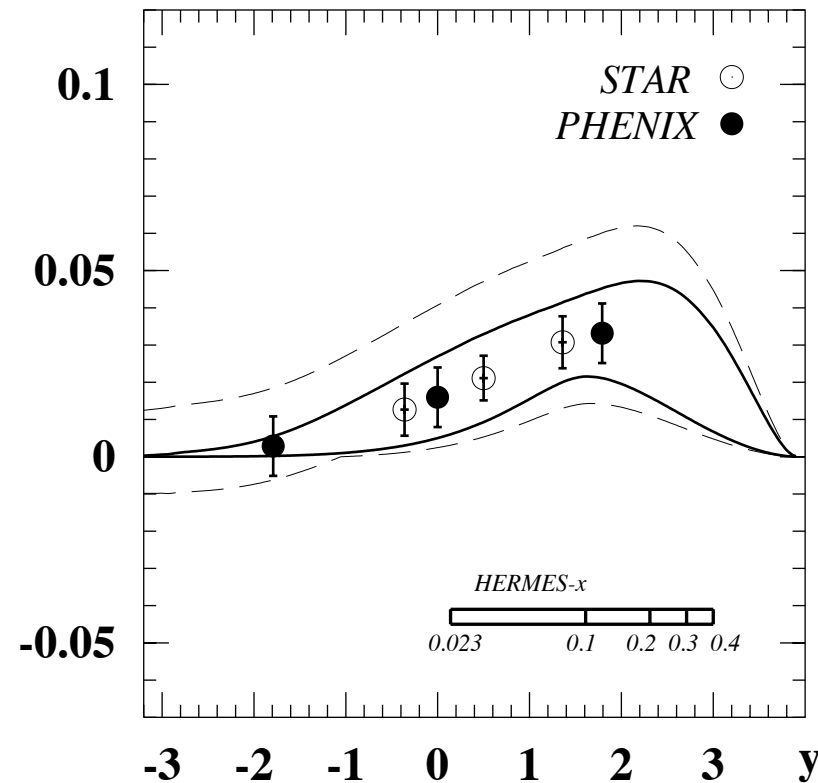


SSA for Drell-Yan dilepton production at RHIC

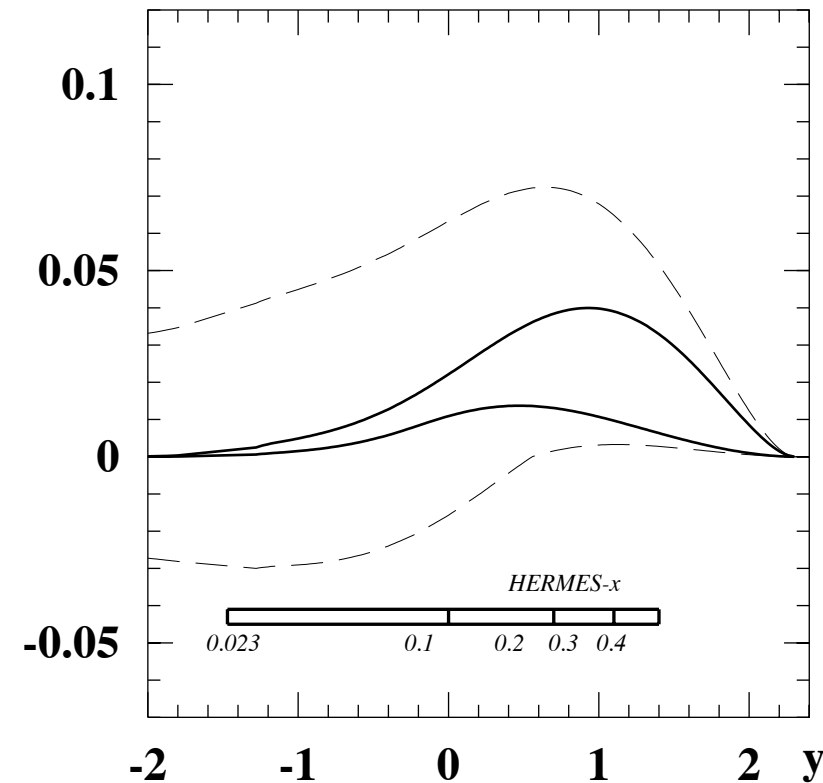
Sum of Sivers functions:

Collins, Efremov, Goeke, Menzel, et.al 2006

$A_{UT}^{\sin(\phi - \phi_S)}$ in $p \uparrow p \rightarrow l^+ l^- X$ at RHIC $Q=4\text{GeV}$



$A_{UT}^{\sin(\phi - \phi_S)}$ in $p \uparrow p \rightarrow l^+ l^- X$ at RHIC $Q=20\text{GeV}$

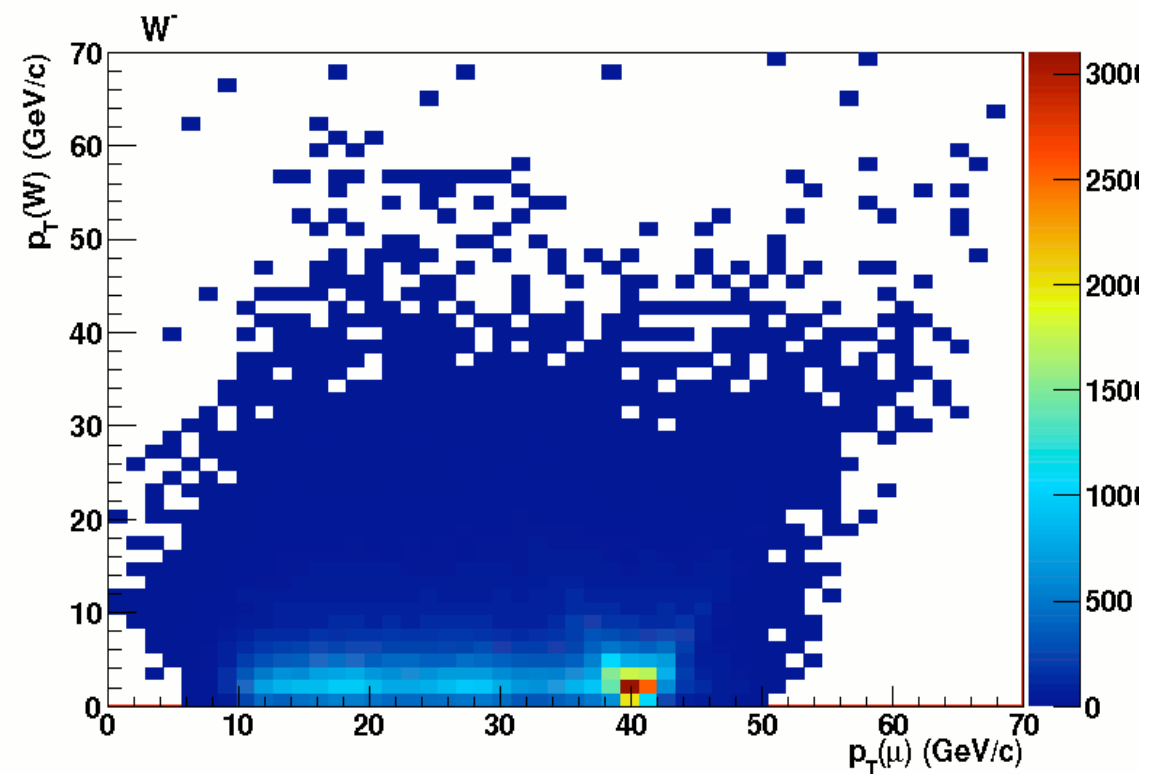
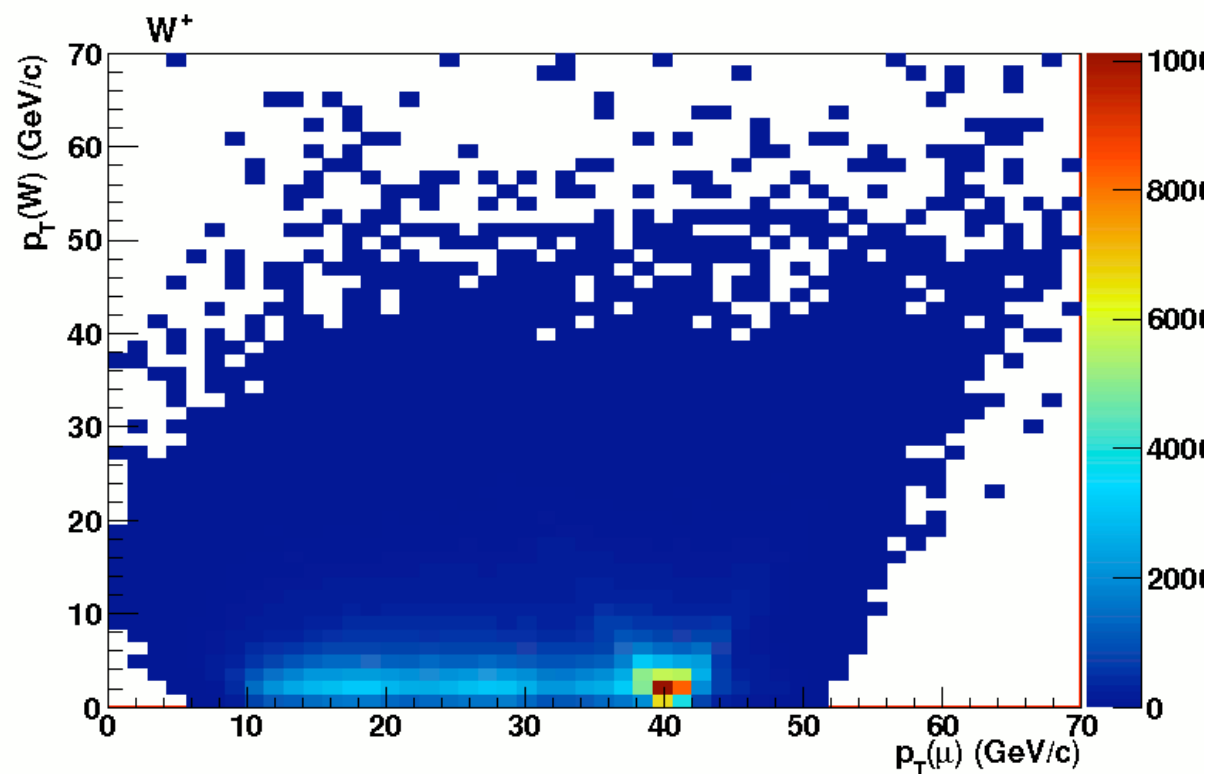


error band: 1- σ uncertainty of the fit of Sivers function

$$A_N = \frac{\sum_q e_q^2 \int \Delta^N f_{q/A\uparrow}(x_1, \mathbf{k}_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})}{2 \sum_q e_q^2 \int f_{q/A}(x_1, k_{\perp 1}) f_{\bar{q}/B}(x_2, k_{\perp 2})} \propto \frac{4}{9} \Delta^N f_u + \frac{1}{9} \Delta^N f_d$$

Production of W^+ , W^- at RHIC

- W boson are primarily produced at the region: $M_W \gg q_T \sim 2\text{GeV}$
 - Lepton $p_T \sim M_W/2$



Courtesy of Kempel, Lajoie (PHENIX)

- TMD approach could be used
 - W events $\sim 10,000$

SSA of W boson

- Collinear factorization will have problems: Sudakov logarithm, resummation, CSS formalism
- Results in TMD approach:

- Spin-dependent:

$$\frac{d\Delta\sigma_{A^\dagger B \rightarrow W}(\vec{S}_\perp)}{dy_W d^2\mathbf{q}_\perp} = \frac{\sigma_0}{2} \sum_{a,b} |V_{ab}|^2 \int d^2\mathbf{k}_{a\perp} d^2\mathbf{k}_{b\perp} \vec{S}_\perp \cdot (\hat{p}_A \times \hat{\mathbf{k}}_{a\perp}) \\ \times \Delta^N f_{a/A^\dagger}^{\text{DY}}(x_a, k_{a\perp}) f_{b/B}(x_b, k_{b\perp}) \delta^2(\mathbf{q}_\perp - \mathbf{k}_{a\perp} - \mathbf{k}_{b\perp})$$

$$x_a = \frac{M_W}{\sqrt{s}} e^{y_W}, \quad x_b = \frac{M_W}{\sqrt{s}} e^{-y_W}$$

$$W^-: \quad ab = d\bar{u}, \bar{u}d, \dots$$

$$W^+: \quad ab = u\bar{d}, \bar{d}u, \dots$$

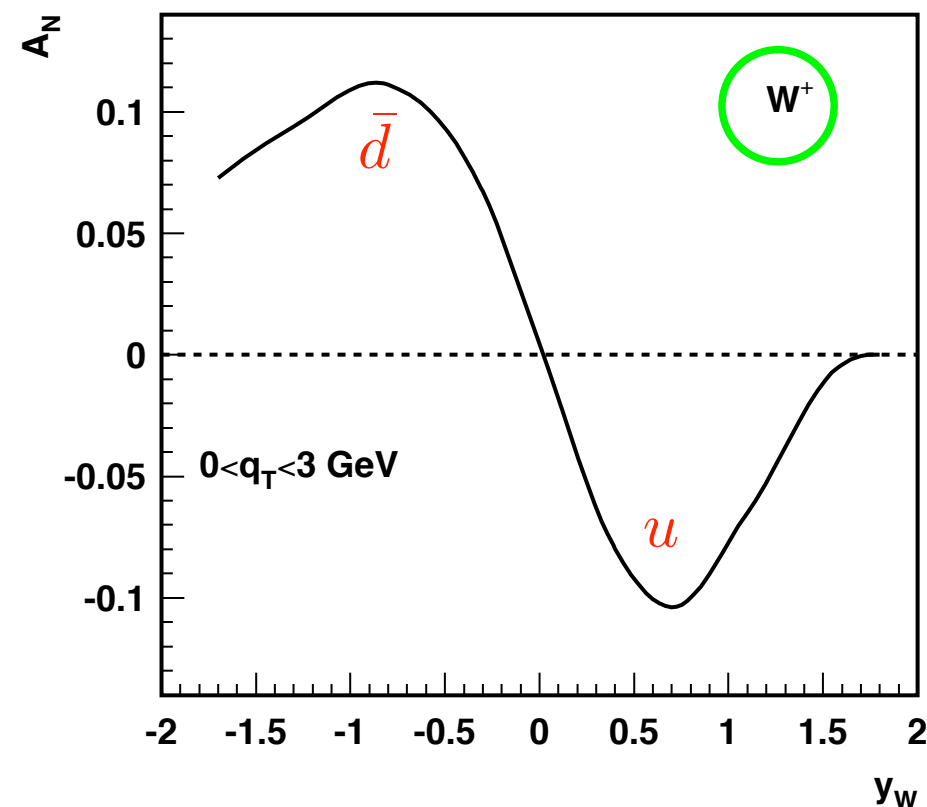
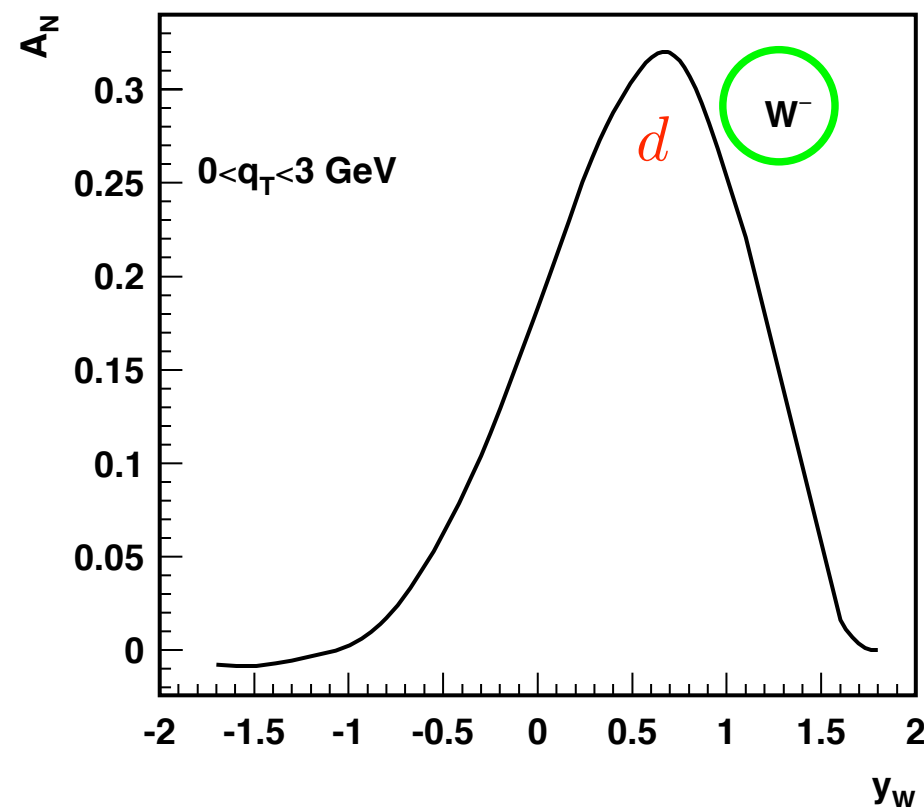
- W- sensitive to d and u-bar Sivers function, W+ sensitive to u and d-bar Sivers function

$$A_N^{(W)} \equiv \frac{d\Delta\sigma(\vec{S}_\perp)_{A^\dagger B \rightarrow W}}{dy_W d^2\mathbf{q}_\perp} \bigg/ \frac{d\sigma_{AB \rightarrow W}}{dy_W d^2\mathbf{q}_\perp}$$

SSA of W bosons: rapidity dependence

Brodsky, Hwang, Schmidt, 2002,
Schmidt, Soffer, 03, Kang, Qiu, 09

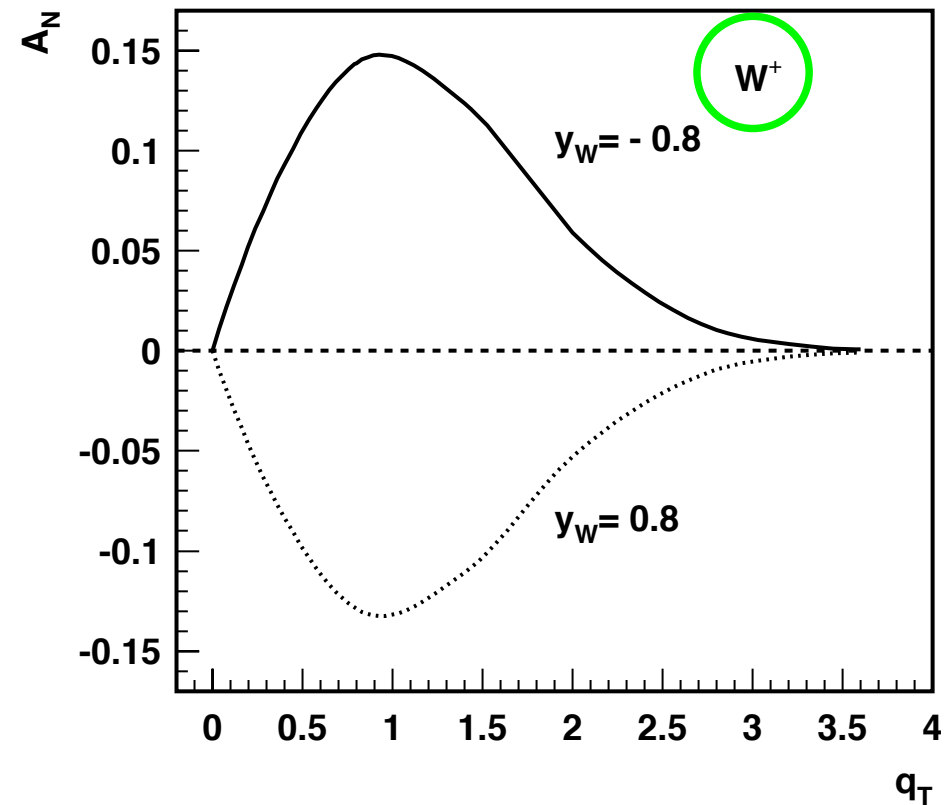
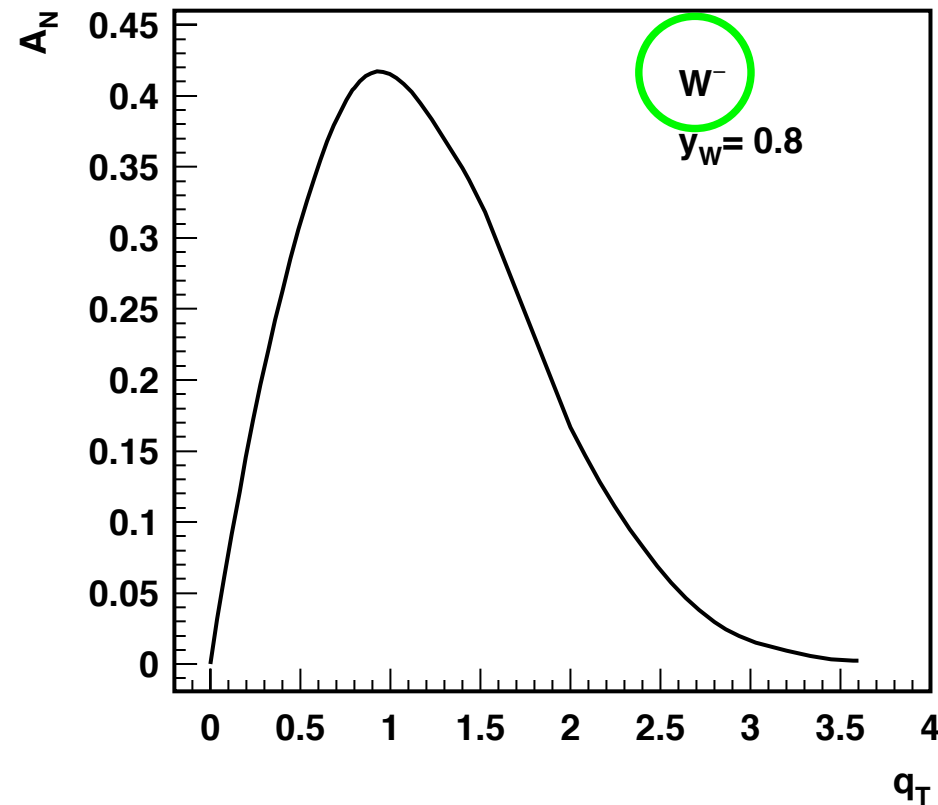
- Use the Sivers functions extracted by Anselmino et al.
- SSAs of W production at RHIC:
 - Sivers function same as DY, different from SIDIS by a sign



- Good flavor separation

SSA of W bosons: q_T dependence

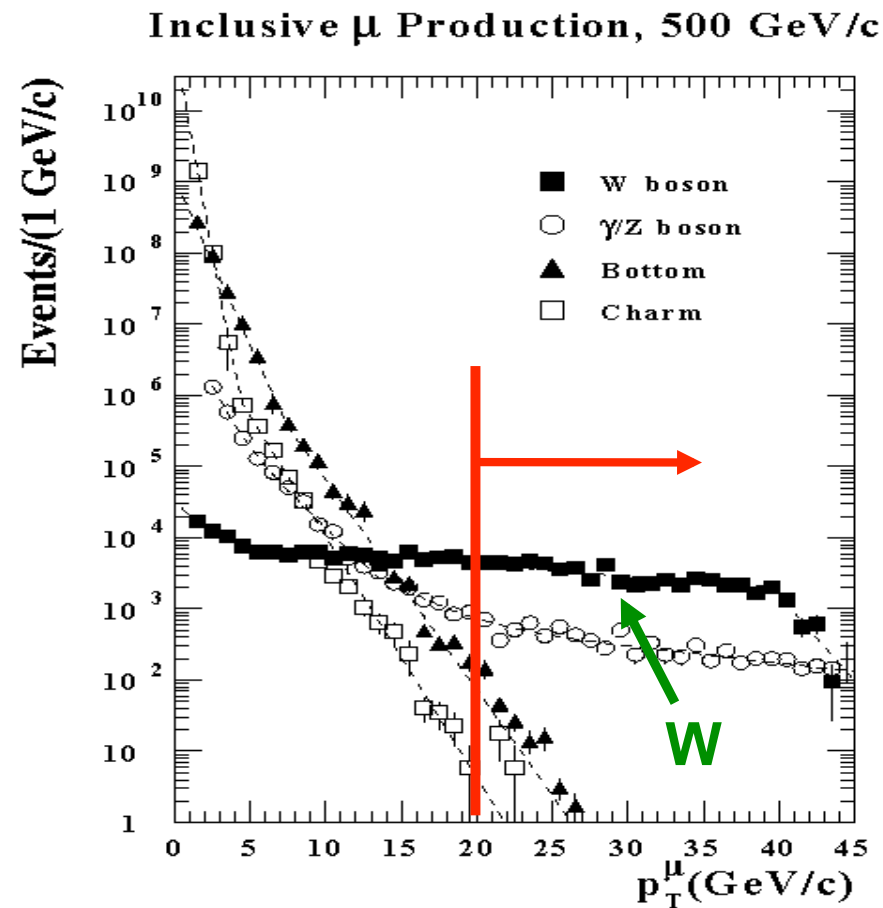
- Large asymmetry: should be able to see sign change



- But, the detectors at RHIC cannot reconstruct the W 's

Lepton from W decay: $W \rightarrow \mu \nu$

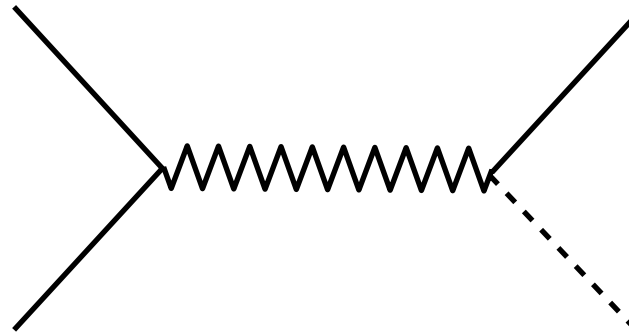
- Inclusive lepton background from Charm/Bottom dies when $p_T > 20$ GeV



- Idea: integrate out the neutrino to measure SSA of inclusive high p_T lepton
- However, $\sin(\Phi_s - \Phi_w)$ dependence of the SSA of W's could dilute the SSA of inclusive lepton

SSA of lepton from W decay: formula

- Spin-dependent:



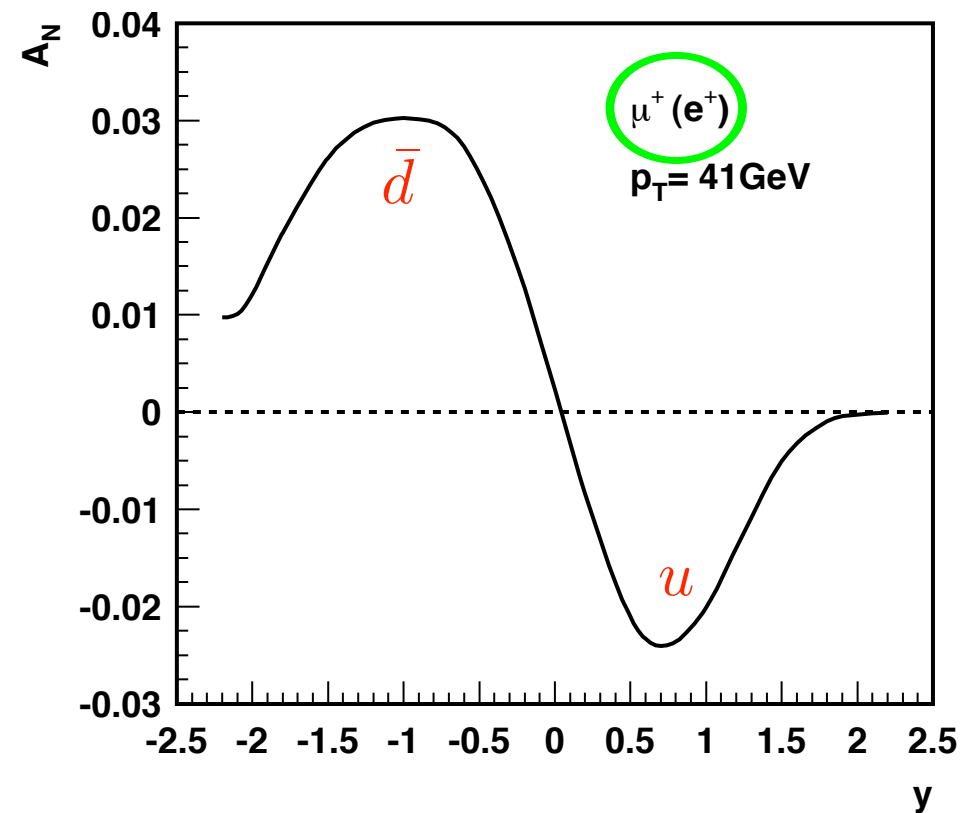
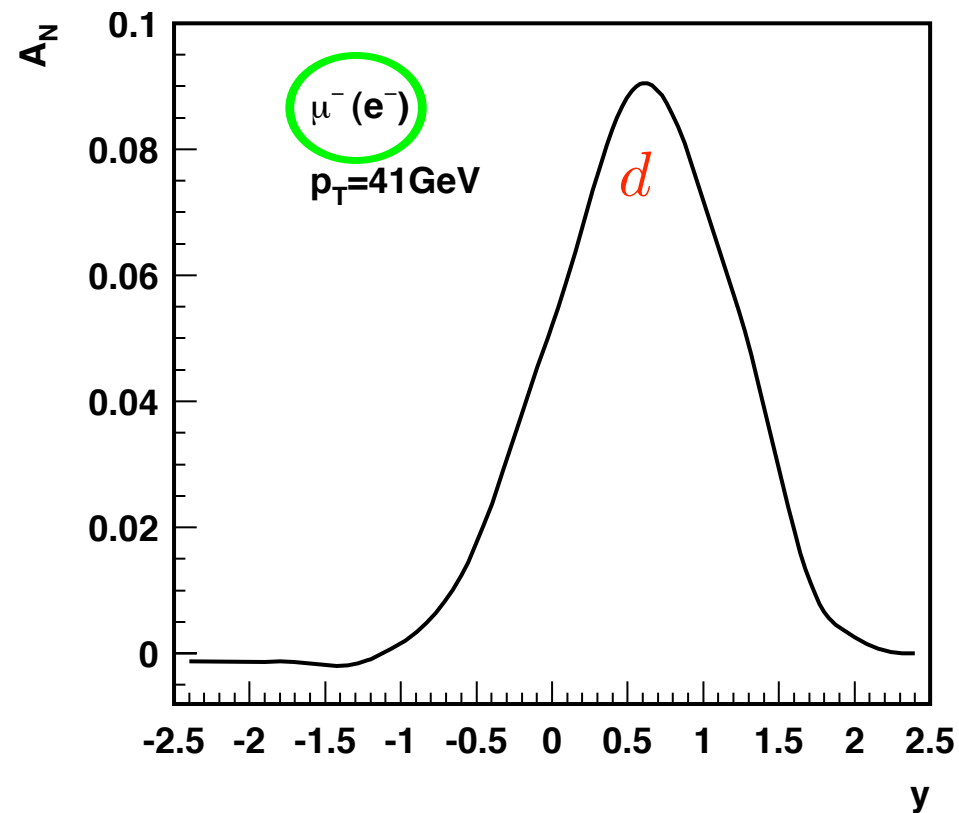
$$\frac{d\Delta\sigma_{A^\uparrow B \rightarrow \ell(p)}(\vec{S}_\perp)}{dy d^2\mathbf{p}_\perp} = \sum_{a,b} |V_{ab}|^2 \int dx_a d^2\mathbf{k}_{a\perp} \int dx_b d^2\mathbf{k}_{b\perp} \vec{S}_\perp \cdot (\hat{p}_A \times \hat{\mathbf{k}}_{a\perp}) \Delta^N f_{a/A^\uparrow}^{\text{DY}}(x_a, k_{a\perp}) \\ \times f_{b/B}(x_b, k_{b\perp}) \frac{1}{16\pi^2 \hat{s}} |\overline{\mathcal{M}}_{ab \rightarrow \ell}|^2 \delta(\hat{s} + \hat{t} + \hat{u})$$

$$|\overline{\mathcal{M}}_{ab \rightarrow \ell}|^2 = \frac{8(G_F M_W^2)^2}{3} \frac{\hat{t}^2}{(\hat{s} - M_W^2)^2 + M_W^2 \Gamma_W^2} \quad ab = \bar{u}d, \bar{u}s, u\bar{d}, u\bar{s}$$

$$|\overline{\mathcal{M}}_{ab \rightarrow \ell}|^2 = \frac{8(G_F M_W^2)^2}{3} \frac{\hat{u}^2}{(\hat{s} - M_W^2)^2 + M_W^2 \Gamma_W^2} \quad ab = d\bar{u}, s\bar{u}, \bar{d}u, \bar{s}u$$

SSA of lepton from W decay: rapidity dependence

- SSA of inclusive lepton is still sufficient for measurement:

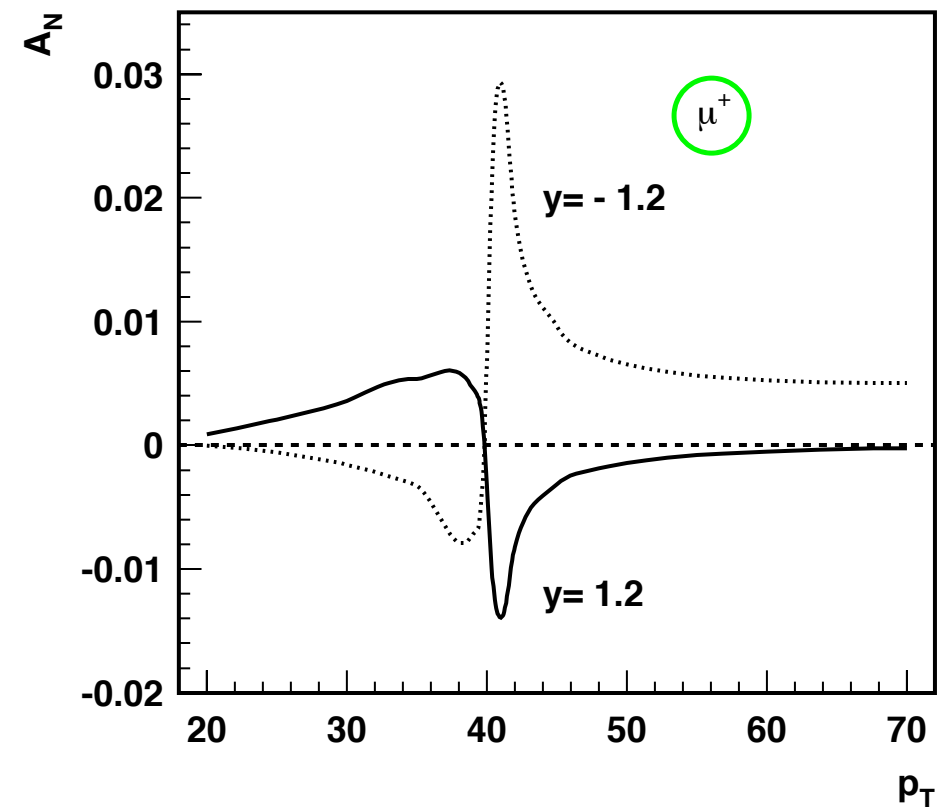
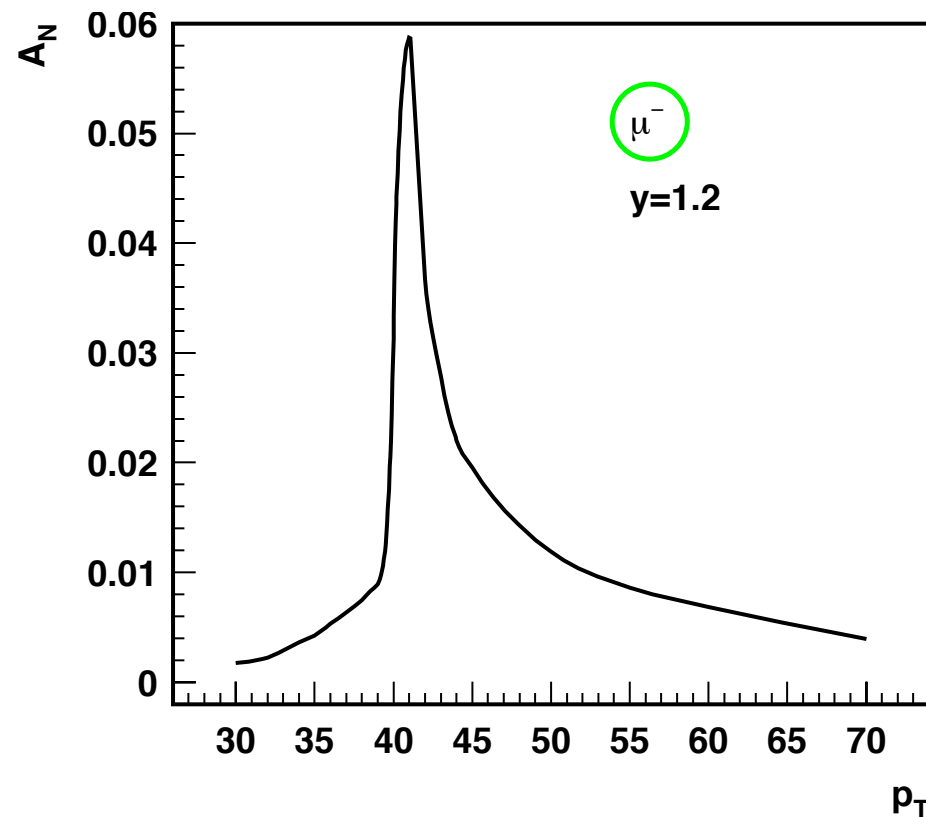


- Good flavor separation:

- $\mu^- (e^-)$ at central-forward rapidity is sensitive to d Sivers function
- $\mu^+ (e^+)$ at forward is sensitive to u Sivers function, at backward is sensitive to \bar{d} Sivers function

SSA of lepton from W decay: p_T dependence

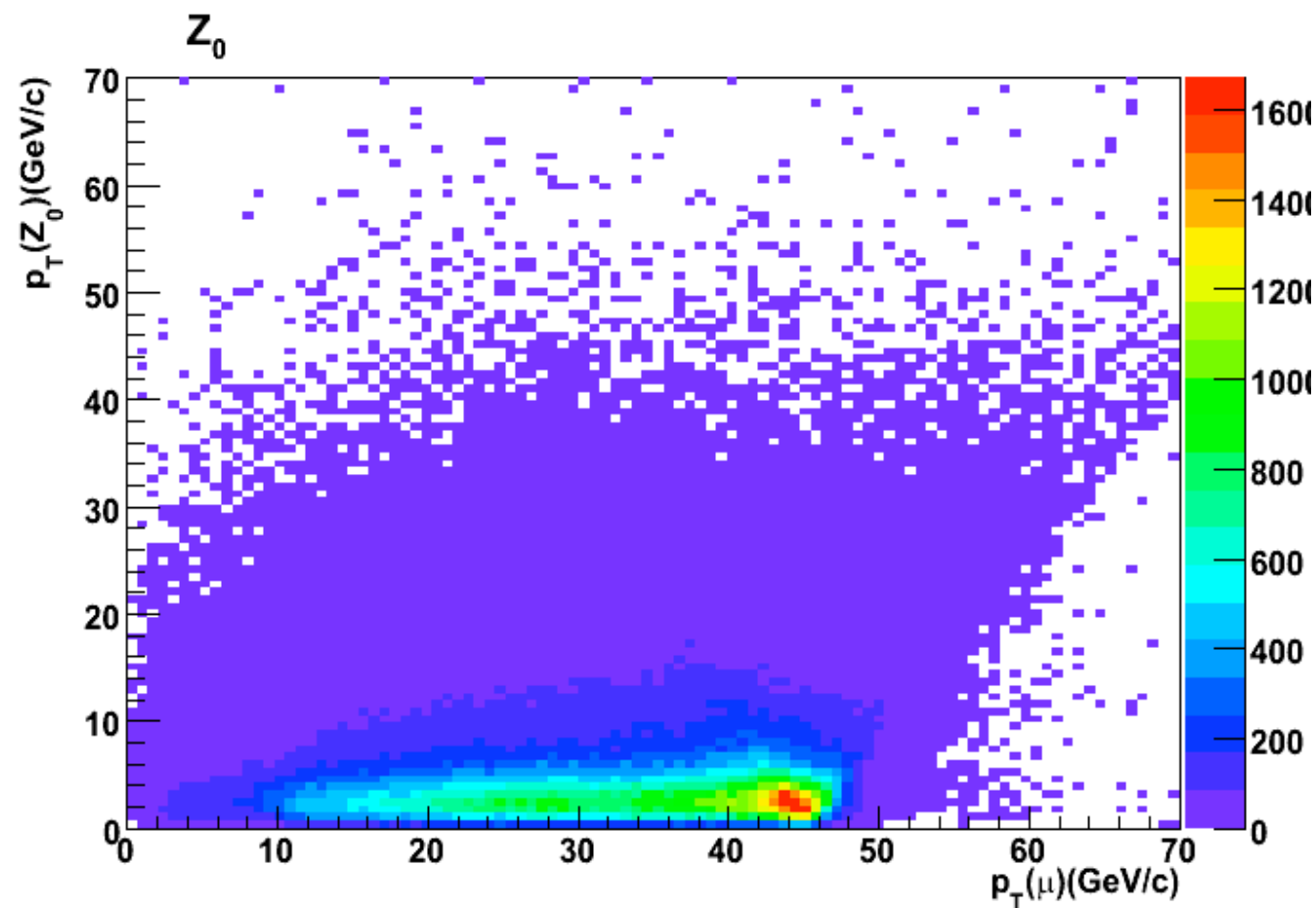
- p_T behavior of SSA of leptons:



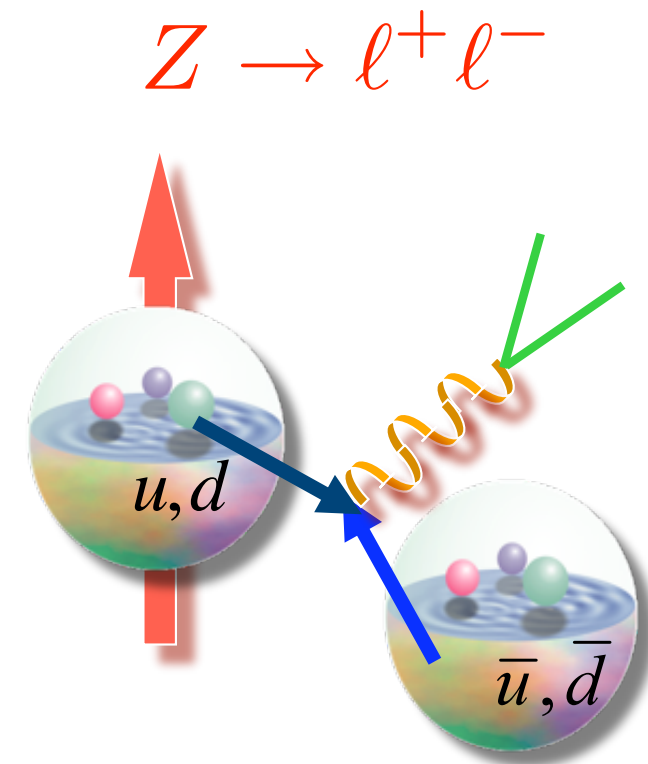
- inherit the key features of W asymmetry
- sharply peaked around $p_T \sim M_W/2$, should help control the potential background

What about Z boson?

- RHIC can reconstruct Z boson



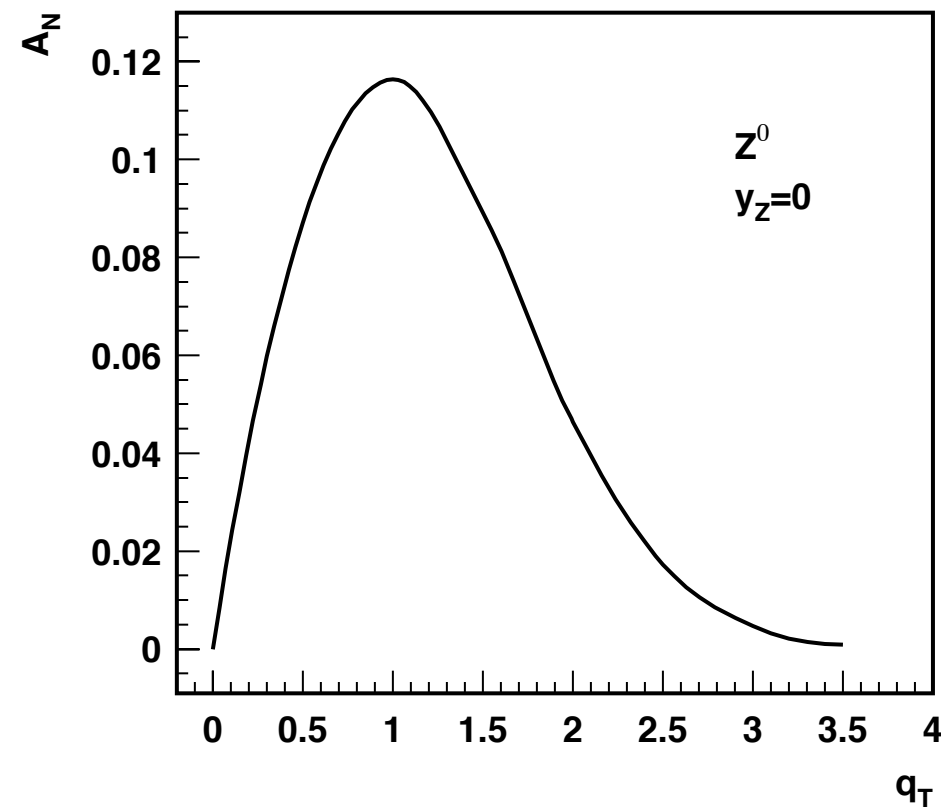
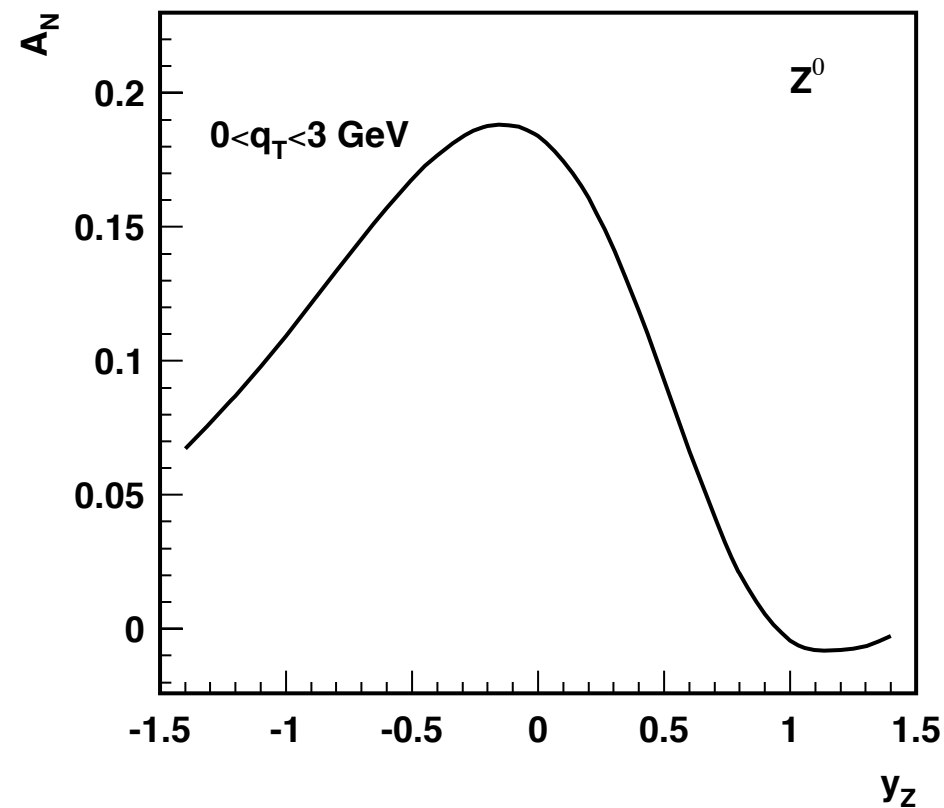
Courtesy of Kempel, Lajoie (PHENIX)



- Events down by an order of magnitude compared to W boson:
 - ~ 1000

SSA of Z boson at RHIC

- Prediction for RHIC kinematics:



- Fairly large asymmetry, should be very good channel to test sign change if one can accumulate enough Z boson



Summary

- Sign change of Sivers function between DY and SIDIS is the most critical test for our current understanding of SSAs
- Besides the standard DY dilepton production, we propose to use the SSAs of W and/or Z boson production to test this sign change
- Lepton decayed from W^+ , W^- boson could give good flavor separation, give even separate tests for Sivers function of different flavors
- Z boson might also be a good channel if one could accumulate enough events

Thank you!